

# Evolution of Supermassive Black Holes Binary Solution

## A. DYNAMICAL FRICTION

**A1.** The deflection angle is defined from:  $\tan \alpha \approx \alpha = \frac{p_y}{p_x}$ , assuming that  $\alpha \ll 1$ . One can find  $p_y = \int F_y dt$ , and according to Newton's gravity law

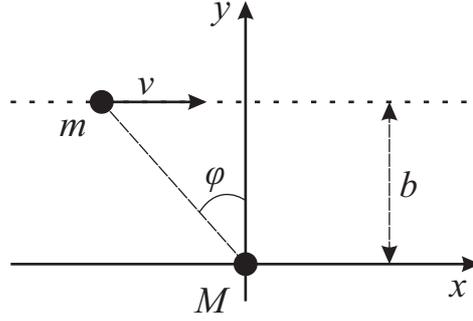
$$F_y = \frac{GMm}{b^2} \cos^3 \varphi$$

The geometry:  $x = b \tan \varphi$ , so we change the variable  $dt = \frac{dx}{v} = \frac{b}{v} \frac{d\varphi}{\cos^2 \varphi}$  and we have

$$p_y = \frac{GMm}{bv} \int_{-\pi/2}^{\pi/2} \cos \varphi d\varphi = \frac{2GMm}{bv}.$$

Here we assume that the body moves along the straight line, due to  $\alpha \ll 1$ , see Fig 1. So  $\alpha = \frac{p_y}{p}$  and

$$\boxed{\alpha = \frac{2GM}{bv^2} = \frac{2b_1}{b}}, \quad \boxed{k = 2}$$



**A2.** During the transit of a massive body, star's energy remains constant:  $p_x^2 + p_y^2 = \text{const.}$  Hence

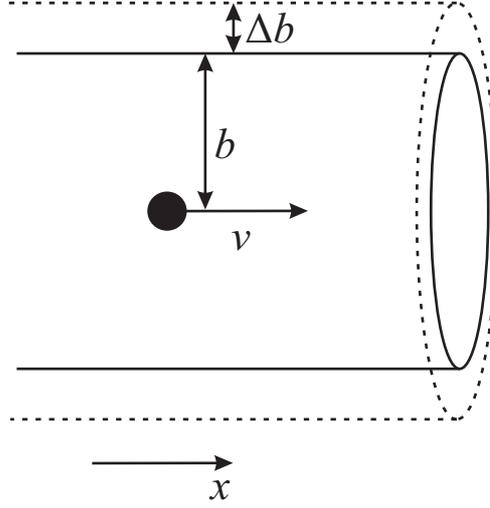
$$(p + \Delta p_x)^2 + p_y^2 = p^2.$$

We know that  $p_y \ll p$ , so the SBH momentum change along the x-axis  $\Delta p_x = -\frac{p_y^2}{2p} = -\frac{\alpha^2}{2}p$ , so

$$\boxed{\Delta p_x = -\frac{2G^2 M^2 m}{b^2 v^3}}.$$

**A3.** To calculate net force we might integrate over stars with different impact parameters. The number of stars' transits during the time  $\Delta t$  equals  $\Delta N = 2\pi b v n db \Delta t$ , so force, decelerating the object along the x-axis,

$$(1) \quad F_{DF} = \frac{1}{\Delta t} \int \Delta p_x dN = -4\pi G^2 M^2 \frac{nm}{v^2} \int_{b_{min}}^{b_{max}} \frac{db}{b} = -4\pi G^2 M^2 \frac{\rho}{v^2} \log \Lambda$$



The above formulas are true only for  $b > b_1$ , so the lower integration limit is  $b_{min} = b_1$ , and the upper limit is determined by the galaxy size  $b_{max} = R$ . So we have

$$(2) \quad F_{DF} = -4\pi G^2 M^2 \frac{\rho}{v^2} \log \Lambda$$

where  $\Lambda = R/b_1$ .

**A4.** We calculate:  $b_1 = \frac{GM}{v^2} = 10.7 \text{ pc}$ ,  $\log \Lambda = 7.5$ .

## B. GRAVITATIONAL SLINGSHOT

**B1.** From the second Newton's law

$$\frac{Mv^2}{a} = \frac{GM^2}{4a^2},$$

and we have for the orbital velocity  $v_{bin} = \sqrt{\frac{GM}{4a}}$ . The system energy is

$$E = E_{kin} + U = 2 \cdot \frac{Mv^2}{2} - \frac{GM^2}{2a}.$$

The answer is

$$(3) \quad \boxed{E = -\frac{GM^2}{4a}}$$

**B2.** From angular momentum conservation law

$$b\sigma = r_m v_0,$$

express  $v_0$ . Write down the energy conservation law

$$\frac{\sigma^2}{2} = \frac{v_0^2}{2} - \frac{GM_2}{r_m}$$

and derive

$$b = r_m \sqrt{1 + \frac{2GM_2}{\sigma^2 r_m}}.$$

**B3.** To estimate the time between collisions let us use an analogy with the gas. As known from the molecular kinetic theory, given that molecules have radii  $r$ , thermal velocities  $v$ , and the molecular concentration  $n$ , the time  $\Delta t$  between collisions of one molecule with the others can be estimated from the relation  $\pi r^2 v n \Delta t = 1$ . In our problem  $b_{max}$  stands in place of the molecule radius, therefore for estimation it can be written

$$(\Delta t)^{-1} = \pi \sigma b_{max}^2 n.$$

Estimate the maximal impact parameter  $b_{max}$ , corresponding to the star collision with the binary system. The star should reach the distance of  $a$  to the binary system to collide. The star at large distances from the SBH binary interacts with it as with a point object of mass  $M_2 = 2M$ . From the results of B.2, assuming  $r_m = a$ , we obtain  $b_{max} = a \sqrt{1 + \frac{4GM}{\sigma^2 a}}$ . Taking into account that  $\sigma^2 \ll \frac{GM}{a}$ , simplify:

$$b_{max} = \frac{2}{\sigma} \sqrt{GMa},$$

so we have

$$\boxed{\Delta t = \frac{m\sigma}{4\pi GM\rho a}}$$

**B4.** During the one act of gravitational slingshot, star energy increases at average by

$$\Delta E_{star} = \frac{mv_{bin}^2}{2} - \frac{m\sigma^2}{2}.$$

So the binary energy decreases by the same magnitude  $\Delta E_{bin} = -\Delta E_{star}$ . Taking into account that  $\sigma \ll v_{bin}$ , we derive

$$\Delta E_{bin} = -\frac{m}{2}v_{bin}^2 = \frac{GmM}{8a}.$$

Average binary system energy loss rate equals

$$(4) \quad \boxed{\frac{dE}{dt} = \frac{\Delta E}{\Delta t} = -\frac{\pi G^2 M^2 \rho}{2\sigma}}$$

Taking the time derivative of (3), we have

$$(5) \quad \frac{dE}{dt} = \frac{d}{dt} \left( -\frac{GM^2}{4a} \right) = \frac{GM^2}{4a^2} \frac{da}{dt},$$

From (4) and (5) the orbit radius variation rate can be estimated as

$$(6) \quad \boxed{\frac{da}{dt} = -\frac{2\pi G\rho a^2}{\sigma}}$$

**B5.** Equation (6) can be easily integrated

$$(7) \quad \frac{da}{a^2} = -\frac{2\pi G\rho}{\sigma} dt.$$

To reduce the radius twice it takes time

$$\boxed{T_{SS} = \frac{\sigma}{2\pi G\rho a_1} = 7.3 \times 10^{-4} \text{ Gy}}$$

### C. EMISSION OF GRAVITATIONAL WAVES

**C1.** Using that  $\omega = \frac{v_{bin}}{a} = \sqrt{\frac{GM}{4a^3}}$  and formulas from the problem text one can obtain:

$$(8) \quad \frac{dE}{dt} = -\frac{1024 \times 4}{5} \times \frac{GM^2 v_{bin}^6}{c^5 a^2} = \frac{64}{5} \cdot \frac{G^4 M^5}{c^5 a^5}.$$

Combining (5) and (8) we get the desirable result:

$$(9) \quad \boxed{\frac{da}{dt} = -\frac{256}{5} \cdot \frac{G^3 M^3}{c^5 a^3}}$$

**C2.** Integrating the equation (9) one can obtain:

$$(10) \quad a^3 da = -\frac{256}{5} \cdot \frac{G^3 M^3}{c^5} dt \quad \implies \quad \frac{a_2^4 - r_g^4}{4} = \frac{256}{5} \cdot \frac{G^3 M^3}{c^5} \cdot T_{GW};$$

And taking into account  $a_2 \gg r_g$  we derive the final result for  $T_{GW}$ :

$$(11) \quad \boxed{T_{GW} = \frac{5}{1024} \cdot \frac{a_2^4 c^5}{G^3 M^3}}$$

**C3.** From the previous equation and  $T_{GW} = t_H$ :

$$(12) \quad \boxed{a_H = \sqrt[4]{\frac{1024}{5} \cdot \frac{G^3 M^3 t_H}{c^5}} = 0.098 \text{ pc}}$$

### D. FULL EVOLUTION

**D1.** The galaxy is spherically symmetric, so mass enclosed within a sphere of radius  $r$  equals

$$(13) \quad m(r) = \int_0^r 4\pi x^2 \rho(x) dx = \frac{\sigma^2 r}{G}.$$

Thus the free fall acceleration of the body equals in the gravitational field of stars is

$$(14) \quad g(r) = \frac{Gm(r)}{r^2} = \frac{\sigma^2}{r}.$$

Therefore the body velocity is determined by relation

$$\frac{v^2}{r} = g = \frac{\sigma^2}{r},$$

which means

$$(15) \quad \boxed{v = \sigma}$$

So the velocity is constant.

**D2.** The energy of SBH in this gravitational field is

$$E = \frac{M\sigma^2}{2} + U$$

So the kinetic energy is constant and

$$\frac{dE}{dt} = \frac{dU}{dt} = \frac{dU}{da} \frac{da}{dt}$$

From the definition of potential energy we have

$$\frac{dU}{da} = g(a)M = \frac{M\sigma^2}{a}$$

Using the result of A3 we have

$$\frac{dE}{dt} = -F_{Df}v = -4\pi G^2 M^2 \frac{\rho(a)}{\sigma} \log \Lambda = -\frac{GM^2 \sigma \log \Lambda}{a^2}.$$

Combining this equations we get the answer

$$(16) \quad \frac{da}{dt} = -\frac{GM \log \Lambda}{a\sigma}$$

**D3.** To estimate one can assume that SBHs form a binary when the mass of stars inside the sphere of radius  $a$  equals to  $M$ :

$$m(a) = \frac{\sigma^2 a}{G} = M,$$

so

$$\boxed{a_1 = \frac{GM}{\sigma^2} = 10.8 \text{pc}}$$

Alternative variant: the force from another SBH is equal to force from all stars:

$$\frac{Gm(a)}{a^2} = \frac{GM}{4a^2}$$

so the answer is

$$\boxed{a_1 = \frac{GM}{4\sigma^2} = 2.7 \text{pc}}$$

**D4.** Integrating the equation (16) we have

$$\frac{a_0^2 - a_1^2}{2} = \frac{GM \log \Lambda}{\sigma} T_1$$

and using that  $a_1 \ll a_0$  we have

$$T_1 = \frac{a_0^2 \sigma}{2GM \log \Lambda} = 0.121 \text{Gy}.$$

**D5.** Total energy losses are caused by gravitational slingshot and gravitational waves emission, so combining equations (4) and (8):

$$(17) \quad \frac{dE}{dt} = -\frac{\pi G^2 M^2 \rho_1}{2\sigma} - \frac{64}{5} \cdot \frac{G^4 M^5}{c^5 a^5}$$

where

$\rho_1 = \rho(a_1) = \rho(10.8\text{pc}) = 6.3 \times 10^3 M_s/\text{pc}^3$ ,      alternative:  $\rho_1 = \rho(2.7\text{pc}) = 1.0 \times 10^5 M_s/\text{pc}^3$   
 Energy losses due to GW dominates when  $\frac{\pi G^2 M^2 \rho_1}{2\sigma} < \frac{64 G^4 M^5}{5 c^5 a^5}$  i.e.  $a < a_2$  where

$$a_2^5 = \frac{128}{5\pi} \cdot \frac{G^2 M^3 \sigma}{c^5 \rho_1} = \frac{512}{5} \cdot \frac{G^3 M^3 a_1^2}{c^5 \sigma}$$

Numerical answer is  $\boxed{a_2 = 0.018 \text{ pc}}$  (alternative:  $a_2 = 0.010 \text{ pc}$ ).

**D6.** For rough approximation it can be considered that at the slingshot stage ( $a > a_2$ ) energy losses are caused only by slingshot, so  $T_2$  is calculated analogously to B5:  $\frac{da}{a^2} = -\frac{2\pi G \rho}{\sigma} dt$  and

$$T_2 \approx \frac{\sigma}{2\pi G \rho_1 a_2} = 0.063 \text{ Gy} \quad (T_2 \approx 0.0068 \text{ Gy})$$

And at the GW emission stage ( $a < a_2$ ) energy losses are caused only by GW emission, so  $T_3$  is calculated directly from C2:

$$T_3 \approx \frac{5}{1024} \cdot \frac{a_2^4 c^5}{G^3 M^3} = \frac{1}{8\pi} \cdot \frac{\sigma}{G \rho_1 a_2} = 0.016 \text{ Gy} \quad (T_3 \approx 0.0017 \text{ Gy})$$

**D7.** Total time of SBH binary evolution from the moment of galaxies merging to SBH merging equals

$$T_{ev} = T_1 + T_2 + T_{GW} = 0.12 + 0.06 + 0.02 \text{ Gy} = 0.20 \text{ Gy} \quad (T_{ev} = 0.13 \text{ Gy})$$